

An analysis of $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ decays at the LHCb experiment

L. Pescatore*,
on behalf of the LHCb collaboration
* *University of Birmingham, United Kingdom*



Abstract

The branching fraction of the rare decay $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ is measured as a function of q^2 , the square of the dimuon invariant mass. The analysis is performed using proton-proton collision data, corresponding to an integrated luminosity of 3.0 fb^{-1} , collected by the LHCb experiment. Evidence of signal is found for the first time in the q^2 region below the square of the J/ψ mass. In the q^2 intervals where the signal is observed, angular distributions are studied and two forward-backward asymmetries, in the dimuon and hadronic systems, are measured for the first time.

1 Rare decays and Λ_b^0

The $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ decay is a rare ($b \rightarrow s$) flavour changing transition, which is forbidden at tree level in the Standard Model (SM) but can happen through loop level electroweak processes¹. Since the branching ratio of this type of decays is small, typically $\sim 10^{-6}$ or less, they are very sensitive to contributions from physics beyond the SM. The study of Λ_b^0 decays is of particular interest for several reasons. First of all, the Λ_b^0 baryon has non-zero spin, which allows to extract additional information about the helicity structure of the underlying theory. Secondly, the Λ_b^0 baryon can be considered as a heavy quark plus a light di-quark system, and therefore the hadronic physics differs significantly from similar B meson decays. Finally, a further motivation specific to the $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ decay is that the Λ baryon decays weakly and its polarisation is preserved which gives access to complementary information to that available from meson decays². In this work the differential branching ratio and angular observables are measured as a function of the square of the dimuon mass, q^2 , since theoretically different treatments of form factors are used depending on the considered q^2 region². Previous observations of the decay^{3 4} $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ found evidence for signal only in q^2 intervals above the square of the mass of the J/ψ resonance. These proceedings are based on an analysis⁵ using pp collision data, corresponding to an integrated luminosity of 3.0 fb^{-1} , collected during 2011 and 2012 by the LHCb detector⁶ at centre-of-mass energies of 7 and 8 TeV, respectively.

2 Differential branching fraction

In order to measure their differential branching fraction, candidate $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ decays are reconstructed from a dimuon and a Λ candidate, where the Λ baryon is reconstructed through its $p\pi^-$ decay mode. The rates are normalised with respect to the tree level $b \rightarrow c\bar{s}$ decay $\Lambda_b^0 \rightarrow J/\psi\Lambda$, where the J/ψ meson decays into a $\mu^+\mu^-$ pair, as this mode has the same final state particles as the signal channel.

The selection is based on a neural network classifier⁷. The signal sample used to train the neural network consists of simulated $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ events, while the background is taken from data in the upper sideband of the Λ_b^0 candidate mass spectrum. The variable that provides the greatest discrimination is the χ^2 from a kinematic fit of the complete decay chain in which the proton and pion are constrained such that the $p\pi^-$ invariant mass corresponds to the known Λ baryon mass, and the Λ and dimuon systems are constrained to originate from their respective vertices. Other variables that contribute significantly are: the transverse momentum of the Λ candidate; the particle identification information for the muons, a likelihood variable built using information from LHCb's muon system and ring imaging Cherenkov detectors⁶; the separation of the muons, the pion and the Λ candidate from the pp interaction vertex; and the distance between the Λ_b^0 and Λ decay vertices.

Since the Λ baryon is long-lived, its decay has a topology with a displaced secondary vertex. This leads to little background from other decays: the only relevant contribution comes B^0 decays into K_s^0 and muons, where $K_s^0 \rightarrow \pi^+\pi^-$ and one pion is misidentified as a proton. This background, especially significant for the normalisation channel, is modelled in the fit.

In Fig. 1 the invariant mass distribution of the $\Lambda\mu^+\mu^-$ system is reported in the q^2 interval $15 < q^2 < 20 \text{ GeV}^2/c^4$ together with the fit function used to extract the yield. The signal is modelled using the sum of two Crystal Ball functions and the combinatorial background with an exponential. The background from K_s^0 decays is very small in the rare decay sample and is not visible on Fig. 1.

The absolute branching fraction of the $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ decay is obtained by multiplying the relative branching fraction by the absolute branching fraction of the normalisation channel⁸, $\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi\Lambda) = (6.3 \pm 1.3) \times 10^{-4}$. Measured values are given in Fig. 1 as a function of q^2 together with SM predictions⁹. The uncertainty on these values is dominated by the precision of the branching fraction for the normalisation channel, while the uncertainty on the relative branching fraction is dominated by the size of the data sample available.

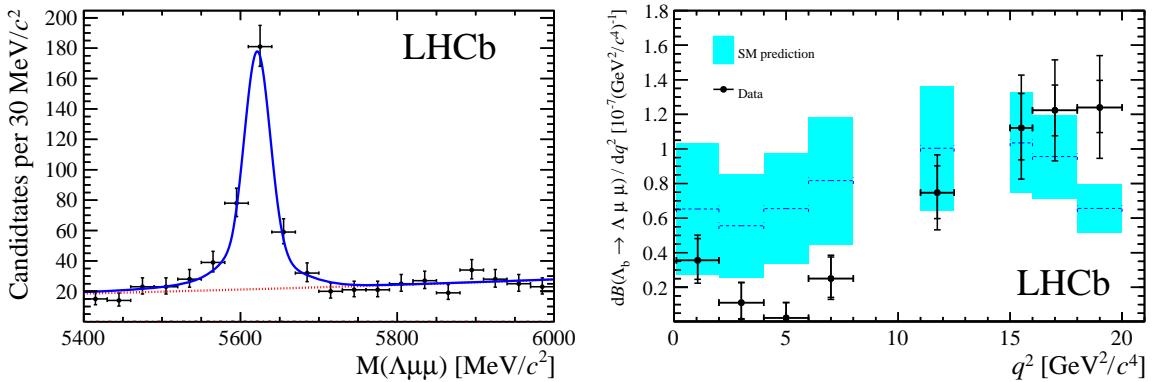


Figure 1 – (left) Invariant mass of the $\Lambda\mu^+\mu^-$ system in the $15 < q^2 < 20 \text{ GeV}^2/c^4$ interval with the fit function overlaid. (right) Measured differential branching fraction of the rare $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ decay as a function of q^2 with the SM prediction⁹ superimposed. The inner error bar on data points represents the total uncertainty on the relative branching fraction (statistical and systematic); the outer error bar also includes the uncertainties from the branching fraction of the normalisation mode.

3 Angular analysis

The $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ decay has a non-trivial angular structure which, in the case of unpolarised Λ_b^0 production, is described by the helicity angles of the muon (θ_ℓ) and proton (θ_h), the angle between the planes defined by the Λ decay products and the two muons, and the square of

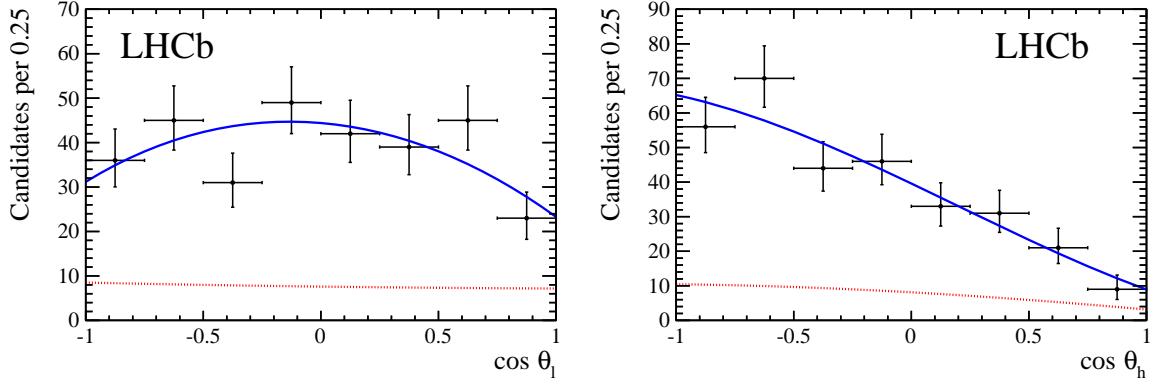


Figure 2 – Distributions of the (left) $\cos \theta_\ell$ and (right) $\cos \theta_h$ angles in data in the $15 < q^2 < 20 \text{ GeV}^2/c^4$ interval with overlaid the total fit function, solid (blue) line, and the background component, dashed (red) line.

the dimuon invariant mass, q^2 . The angle θ_ℓ is defined as the angle between the positive (negative) muon and the dimuon system directions and θ_h as the angle between the proton and the Λ baryon directions, both in the Λ_b^0 ($\bar{\Lambda}_b^0$) rest frame.

In this work two forward-backward asymmetries, in the dimuon (A_{FB}^ℓ) and $p\pi^-$ (A_{FB}^h) systems, are measured. The observables are determined from one-dimensional angular distributions as a function of $\cos \theta_\ell$ and $\cos \theta_h$. The differential rate as a function of $\cos \theta_\ell$ is described by the function

$$\frac{d^2\Gamma(\Lambda_b \rightarrow \Lambda \ell^+ \ell^-)}{dq^2 d\cos \theta_\ell} = \frac{d\Gamma}{dq^2} \left[\frac{3}{8} (1 + \cos^2 \theta_\ell) (1 - f_L) + A_{\text{FB}}^\ell \cos \theta_\ell + \frac{3}{4} f_L \sin^2 \theta_\ell \right], \quad (1)$$

where f_L is the fraction of longitudinally polarised dimuons. The rate as a function of $\cos \theta_h$ has the form

$$\frac{d^2\Gamma(\Lambda_b \rightarrow \Lambda(\rightarrow p\pi^-) \ell^+ \ell^-)}{dq^2 d\cos \theta_h} = \mathcal{B}(\Lambda \rightarrow p\pi^-) \frac{d\Gamma(\Lambda_b \rightarrow \Lambda \ell^+ \ell^-)}{dq^2} \frac{1}{2} (1 + 2A_{\text{FB}}^h \cos \theta_h). \quad (2)$$

These expressions assume that Λ_b^0 baryons are produced unpolarised, which is in agreement with the Λ_b^0 production polarisation recently measured at LHCb¹⁰.

The observables are measured using unbinned maximum likelihood fits. The signal PDF consists of a theoretical shape, given by Eqs. 1 and 2, multiplied by a function modelling the angular efficiency. Selection requirements on the minimum momentum of the muons generate distortions in the $\cos \theta_\ell$ distribution by removing candidates with extreme values of $\cos \theta_\ell$. Similarly, impact parameter requirements remove events with large $|\cos \theta_h|$, since very forward hadrons tend to have smaller impact parameter values. The angular efficiency is parametrised using a second-order polynomial and determined by fitting simulated events.

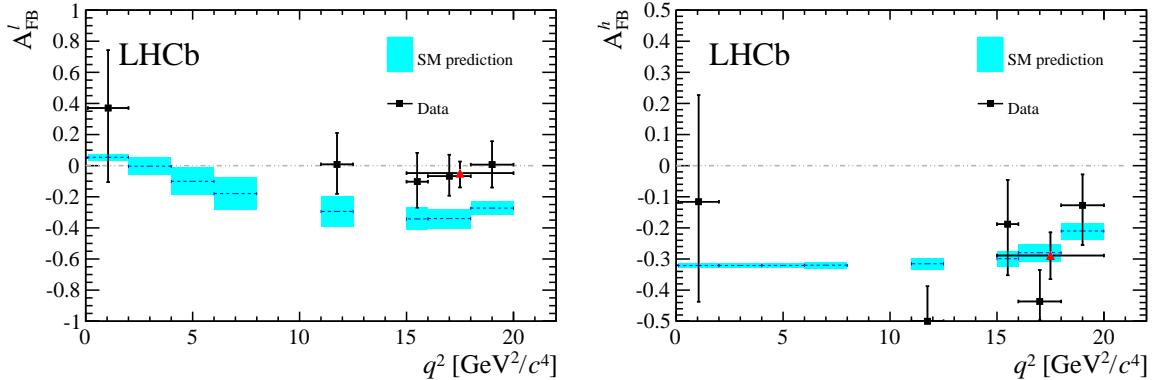


Figure 3 – Forward-backward asymmetries as a function of q^2 in the dimuon system (left) and $p\pi^-$ system (right) with the SM prediction⁹ overlaid.

To limit systematic uncertainties related to the background parametrisation, a narrow interval around the mass peak, dominated by the signal, is used in the angular analysis and a polynomial component is added to the fit to account for the residual background.

Figure 2 shows the $\cos \theta_\ell$ and $\cos \theta_h$ distributions in the $15 < q^2 < 20 \text{ GeV}^2/c^4$ interval with the fit functions overlaid. Measured values of the leptonic and hadronic forward-backward asymmetries, A_{FB}^ℓ and A_{FB}^h , are shown in Fig. 3 together with SM predictions⁹. The statistical uncertainties are obtained using the likelihood-ratio ordering method¹¹ and nuisance parameters are accounted for using the plug-in method¹². One dimensional 68% Confidence Level (CL) intervals are obtained by varying one parameter at a time and treating the others as nuisance parameters. In the analysis⁵ the statistical uncertainties on A_{FB}^ℓ and f_L are also reported in the form of two-dimensional 68% CL regions, where the likelihood-ratio ordering method is applied by varying both observables at the same time and therefore taking correlations into account.

4 Conclusions

A measurement of the differential branching fraction of the rare $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ decay is performed using data recorded by the LHCb detector at centre-of-mass energies of 7 and 8 TeV and corresponding to an integrated luminosity of 3.0 fb^{-1} . Evidence for the signal is found for the first time in the q^2 region below the square of the J/ψ mass, and in particular in the $0.1 < q^2 < 2.0 \text{ GeV}^2/c^4$ interval, where an enhanced yield is expected due to the proximity of the photon pole. The uncertainties of the measurements in the $15 < q^2 < 20 \text{ GeV}^2/c^4$ interval are reduced by a factor of approximately three relative to the previous LHCb measurement³. This improvement is due to a larger data sample and a better control of systematic uncertainties. The branching fraction measurements are compatible with SM predictions in the high- q^2 region, above the square of the J/ψ mass, and lie below the predictions in the low- q^2 region. Furthermore, the first measurement of angular observables for this decay is reported. Two forward-backward asymmetries, in the dimuon and $p\pi^-$ systems, are measured. The measurements of the A_{FB}^h observable are in good agreement with the SM predictions while for the A_{FB}^ℓ observable measurements are consistently above the predictions.

References

1. T. Gutsche *et al.*, *Rare baryon decays $\Lambda_b^0 \rightarrow \Lambda \ell^+ \ell$ ($\ell = e, \mu, \tau$) and $\Lambda_b^0 \rightarrow \Lambda \gamma$: differential and total rates, lepton- and hadron-side forward-backward asymmetries*, *Phys. Rev. D* **87**, 074031 (2013).
2. P. Ber, T. Feldmann, and D. van Dyk, *Angular analysis of the decay $\Lambda_b \rightarrow \Lambda(\rightarrow N\pi)\ell^+\ell^-$* , *JHEP* **01**, 155 (2015).
3. LHCb collaboration, R. Aaij *et al.*, *Measurement of the differential branching fraction of the decay $\Lambda_b^0 \rightarrow \Lambda^0 \mu^+ \mu^-$* , *Phys. Lett. B* **725**, 25 (2013).
4. CDF collaboration, T. Aaltonen *et al.*, *Observation of the baryonic flavor-changing neutral current decay $\Lambda_b^0 \rightarrow \Lambda^0 \mu^+ \mu^-$* , *Phys. Rev. Lett.* **107**, 201802 (2011).
5. LHCb collaboration, R. Aaij *et al.*, *Differential branching fraction and angular analysis of $\Lambda_b^0 \rightarrow \Lambda^0 \mu^+ \mu^-$ decays*, arXiv:1503.07138.
6. LHCb collaboration, A. A. Alves Jr. *et al.*, *The LHCb detector at the LHC*, *JINST* **3**, S08005 (2008).
7. M. Feindt and U. Kerzel, *The NeuroBayes neural network package*, *Nucl. Instrum. Methods A* **559**, 190 (2006).
8. Particle Data Group, K. A. Olive *et al.*, *Review of particle physics*, *Chin. Phys. C* **38**, 090001 (2014).
9. S. Meinel, *Flavor physics with Λ_b^0 baryons*, PoS **LATTICE2013**, 024 (2014).
10. LHCb collaboration, R. Aaij *et al.*, *Measurements of the $\Lambda_b^0 \rightarrow J/\psi \Lambda^0$ decay amplitudes and the Λ_b^0 polarisation in pp collisions at $\sqrt{s} = 7 \text{ TeV}$* , *Phys. Lett. B* **724**, 27 (2013).
11. G. J. Feldman and R. D. Cousins, *Unified approach to the classical statistical analysis of small signals*, *Phys. Rev. D* **57**, 3873 (1998).
12. B. Sen, M. Walker, and M. Woodroffe, *On the unified method with nuisance parameters*, *Statistica Sinica* **19**, 301 (2009).